Comment on “Experimental and Theoretical Constraints of Bipolaronic Superconductivity in High Tc Materials: An Impossibility”

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Additional Information:

- This article was published in the journal, Physical Review Letters [© American Physical Society]. It is also available at: http://link.aps.org/abstract/PRL/v82/p2620.

Metadata Record: https://dspace.lboro.ac.uk/2134/1735

Publisher: © American Physical Society

Please cite the published version.
Comment on “Experimental and Theoretical Constraints of Bipolaronic Superconductivity in High $T_c$ Materials: An Impossibility”

Recently Chakraverty et al. [1] raised objections to the bipolaron model of superconducting cuprates [2]. They claimed that (a) Bose-Einstein condensation (BEC) is not realized for any preformed pairs because this scenario cannot account for the observed value of $T_c$, (b) bipolarons are ruled out as superconducting pairs because they are too heavy, (c) the coherence length in cuprates is such that pairs are strongly overlapped, and (d) the bipolaron theory is incompatible with photoemission experiments. Here I show that their objections are the result of an incorrect approximation for the bipolaron energy spectrum and misuse of our theory.

It is commonly accepted that new superconductors lie close to BEC. We can now assess how close. The bipolaron energy spectrum has been derived for perovskites [3]. It consists of two energy bands $E_k = 2t \cos(k_x a) + t' \cos(k_y a) + t_\perp \cos(k_d)$, where $a$ and $d$ are the lattice constants, and $t, t'$ (not $t/4$ and $t_\perp$) the bipolaron hopping integrals. By expressing the band-structure parameters through the in-plane and out-of-plane penetration depths $\lambda_{ab} = [m_e c^2/8\pi n_B e^2(m_s + m_y)]^{1/2}$, $\lambda_\perp = [m_e c^2/16\pi n_B e^2]^{1/2}$, one readily obtains $T_c$ from the density sum rule [2],

$$k_B T_c = \frac{\hbar^2}{20\lambda_{ab}^2 e^2[1 + \ln(16\pi n_B k_B T c^2 \lambda_\perp^2 d^2 / \hbar^2 c^2)]}.$$  

(1)

Here $m_s = \hbar^2 / 2m_s^*, m_y = 4m_s^*$, and $m_c = \hbar^2 / 2t_\perp d^2$. This expression provides an unambiguous parameter-free estimate of $T_c$. In particular, we obtain $T_c = 100$ K for YBa$_2$Cu$_3$O$_{y}$ ($\lambda_{ab} = 1600$ Å [1], $\lambda_\perp = 12600$ Å [4], and the bipolaron density $n_B = 6 \times 10^{21}$ cm$^{-3}$ assuming that all 2x holes are bound into bipolarons in the doped Mott insulator YBa$_2$Cu$_3$O$_{y+x}$. The measured value is $T_c = 92$ K [4]. Hence, the consideration, which takes into account the multiband energy structure clearly indicates that cuprates are in the BEC regime, while the erroneous approximation [1] yields $T_c$ about 3 times higher.

We now check to what extent our expression for the bipolaron hopping integral, $t = 0.25 t_{\text{band}} e^{-\gamma^2} / \sqrt{2}$ [3] corresponds to the value of the penetration depth. Here $g^2 = \gamma \alpha^2 + \gamma = \sum_q \gamma^2(q) [1 - \cos(q, a)] / \sum_q \gamma^2(q)$ and $\alpha^2 = e_p / \hbar \omega$. A strong interaction of carriers with $c$-axis longitudinal phonons (the frequency $\hbar \omega = 74$ meV) has been established experimentally in YBa$_2$Cu$_3$O$_{y+x}$ while the coupling with other phonons is relatively weak [5]. Because of a low $c$-axis conductivity this interaction is the Fröhlich unscreened interaction with a $q$-dependent matrix element $\gamma(q) \sim g_c / q^2$ [3], which yields $\gamma = 0.162$. As a result, one obtains $m_s = 12 m_e$ with the same polaron level shift $e_p = 250$ meV and bare hopping $t_{\text{band}} = 150$ meV as in Ref. [1]. This bipolaron mass is close to $m_s^* = 14 m_e$ in YBa$_2$Cu$_3$O$_{y}$ obtained from the experimental $\lambda_{ab}$. It is also close to $m^* = 15 m_c$ found in those optical experiments [6], which distinguish between incoherent and Drude contributions. Mott and I always emphasized that small bipolarons in cuprates are inter-site pairs [2] as a result of unscreened interaction [3]. By considering on-site bipolarons or by leaving out the coefficient $\gamma$ in our expression for inter-site bipolaron hopping Chakraverty et al. misused the theory and overestimate the small bipolaron mass by 2 orders of magnitude.

The coherence length of the charged Bose gas has nothing to do with the size of a boson. It can be as large as in the BCS superconductor [2]. Hence, one cannot distinguish the BCS and BEC by arguing (c) that the coherence area is large enough to accommodate many holes. What is really conclusive is the critical behavior. All experiments so far unambiguously show that the upper critical field and the specific heat near the transition are those of charged bosons [2]. The claim (d) that with bipolarons “one should expect essentially $k$-independent spectral functions showing a $k$-independent gap” [1] is not correct either. When an electron is emitted from the sample, the resulting hole propagates in the polaron band. It is the dispersion of this band which is measured by angle-resolved photoemission spectroscopy. The normal state gap is also $k$ dependent owing to the polaron and bipolaron band dispersion. In fact, the bipolaron theory describes all major features of the tunneling and photoemission spectra [7].

To the best of our knowledge there are no unambiguous experimental facts so far which are qualitatively in disagreement with the bipolaron theory. Of course, new facts may disagree. What is clear, however, is that any theory, beautiful or not, cannot be destroyed by “ugly” artifacts as those in Ref. [1].

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Received 17 July 1998 [S0031-9007(99)08745-1]
PACS numbers: 74.20.Mn, 71.38.+i