Microstructure and electrical properties of Aurivillius phase \((\text{CaBi}_2\text{Nb}_2\text{O}_9)_{1-x}(\text{BaBi}_2\text{Nb}_2\text{O}_9)_x\) solid solution

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Microstructure and electrical properties of Aurivillius phase (CaBi$_2$Nb$_2$O$_9$)$_{1-x}$(BaBi$_2$Nb$_2$O$_9$)$_x$ solid solution

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The microstructures and electrical properties of Aurivillius phase ferroelectric solid solutions of (CaBi$_2$Nb$_2$O$_9$)$_{1-x}$(BaBi$_2$Nb$_2$O$_9$)$_x$ ($0 \leq x \leq 1$) have been studied. X-ray diffraction analyses revealed a bismuth layered structure for all compositions. Scanning electron microscope images showed randomly oriented and platelike grain morphology. The Curie point $T_c$ or the maximum permittivity temperature $T_m$ decreased with increasing $x$. The (CaBi$_2$Nb$_2$O$_9$)$_{1-x}$(BaBi$_2$Nb$_2$O$_9$)$_x$ ceramics exhibited a ferroelectric–paraelectric phase transition at small $x$ values ($x \leq 0.5$), whereas a relaxor behavior was observed at high $x$ values ($x \geq 0.8$). The $d_{33}$ value of CaBi$_2$Nb$_2$O$_9$ ceramics was enhanced by Ba$^{2+}$ doping on the A-sites ($x \leq 0.3$). A combination of high $d_{33}$ values and high $T_c$ points ($>700 ^\circ C$) suggests that compositions with $x = 0.3$ could be good candidates for high-temperature piezoelectric applications. The composition with $x = 0.8$ is a relaxor ferroelectric with $T_m$ around 320 $^\circ C$ at 1 MHz. © 2010 American Institute of Physics. [doi:10.1063/1.3457229]

I. INTRODUCTION

Bismuth layer-structured ferroelectrics (BLSFs), with the general formula (Bi$_2$O$_2$)$_{2+}$(A$_{m-1}$B$_2$O$_{3m+1}$)$_{2-}$, are an important family of lead-free ferroelectrics first identified by Aurivillius. The structure of these compounds can be described as pseudo-perovskite (A$_{m-1}$B$_2$O$_{3m+1}$)$_{2-}$ slabs separated by (Bi$_2$O$_2$)$_{2+}$ layers along the crystallographic c-axis. The 12-coordinated A site can be occupied by monovalent, divalent, or trivalent metallic cations. While the octahedral-coordinated B site can be occupied by tetravalent, pentavalent, or hexavalent metallic cations. The number of octahedra along the c-axis between two neighboring (Bi$_2$O$_2$)$_{2+}$ layers is indicated by $m$. Aurivillius phase materials have generated increasing attention due to their potential use in nonvolatile ferroelectric random-access memory$^{3,4}$ and high-temperature piezoelectric applications$^5$ because of their fatigue-free properties and high Curie point, respectively. Some Aurivillius phase compounds show interesting relaxor and multiferroic properties when Ba/lanthanides$^{6,7}$ and Fe (Ref. 8) are on the A- and B-site in the general formula, respectively.

The origin of ferroelectricity in BLSFs is the displacement of the A site cations of the perovskite block with cooperative tilting of the BO$_6$ octahedra.$^9,10$ Compared with traditional piezoelectric ceramics, such as lead zirconate titanate and BaTiO$_3$, BLSFs compounds are characterized by low piezoelectric activity and electromechanical coupling factors and high coercive field because the rotation of the spontaneous polarization is restricted to the a-b plane.$^{11}$ It has been reported that BLSFs solid solution ceramics exhibit improved electrical properties compared to end-member compounds. For example, the mechanical quality factor $Q_m$ reached a maximum value of 11000 at $x = 0.75$ in solid solution of (SrBi$_2$Ta$_2$O$_9$)$_{1-x}$(CaBi$_2$Ta$_2$O$_9$)$_x$.$^{12}$ The $Q_m$ and electromechanical coupling factors $k_p$ were enhanced up to the maximum value of 13500 and 0.12, respectively, in SrBi$_2$Ta$_2$O$_9$–Bi$_3$TiTaO$_9$ solid solution system, Sr$_{0.5}$Bi$_2$Ta$_2$O$_9$–Sr$_{0.5}$Bi$_2$Ta$_2$O$_9$ (1 $\leq x \leq 2$) at $x = 1.25$. The maximum $k_p$ of 9.0% and $Q_m$ of 7000 were obtained for $x = 0.75$ in solid solution system (Na$_{0.5}$Bi$_2$Ta$_2$O$_9$)$_{1-x}$(Sr$_{0.5}$Bi$_2$Ta$_2$O$_9$)$_x$.$^{13}$

Solid solution of (SrBi$_2$Ta$_2$O$_9$)$_{1-x}$(Bi$_3$TiNbO$_9$)$_x$ exhibited at least one order lower conductivity in the temperature range of 400–900 $^\circ C$ than those of the end members SrBi$_2$Ta$_2$O$_9$ and Bi$_3$TiNbO$_9$ and remanant polarization ($P_r$) of (SrBi$_2$Ta$_2$O$_9$)$_{1-x}$(Bi$_3$TiNbO$_9$)$_x$ with $x = 0.8$ was 14.4 $\mu$C/cm$^2$, which was higher than $P_r$ (=11.8 $\mu$C/cm$^2$) of SrBi$_2$Ta$_2$O$_9$.$^{14}$ (Bi$_3$TiNbO$_9$)/(SrBi$_2$Ta$_2$O$_9$)$_{1-x}$ solid solution ceramics with $x = 0.6$ present a higher piezoelectric coefficient ($d_{33}$ = 11 pC/N) and lower coercive field ($E_c$ = 0.67 kV/cm) than those of Bi$_3$TiNbO$_9$ ($d_{33}$ = 1–4 pC/N, $E_c$ > 75 kV/cm).$^{15}$ In (Bi$_3$TiNbO$_9$)$_{1-x}$(BaBi$_2$Nb$_2$O$_9$)$_x$, much larger $P_r$ value (6–7 $\mu$C/cm$^2$) was observed for $x = 0.1–0.3$ compared with $P_r$ (1 $\mu$C/cm$^2$) of Bi$_3$TiNbO$_9$. The resistivity for $x = 0.2$ was one order larger than that of Bi$_3$TiNbO$_9$.$^{16}$

In addition, compared to normal ferroelectrics, relaxor ferroelectrics (RFE) have attracted intensive interest due to their large dielectric permittivity and electromechanical constants for piezoelectric transducer and actuator applications.$^{18}$ But traditional RFE piezoelectrics, such as Pb(Mg$_{1/3}$Nb$_{2/3}$)O$_3$–PbTiO$_3$ and Pb(Zn$_{1/3}$Nb$_{2/3}$)O$_3$–PbTiO$_3$, are not suitable for application under harsh thermal conditions due to their relatively low $T_m$ temperatures (<300 $^\circ C$), at which there is the maximum of dielectric permittivity.$^{19,20}$ There are relatively few reports on relaxor behavior in BLSFs solid solution ceramics. The study of (Na$_{0.5}$Bi$_{1.5}$Ti$_3$O$_{15}$)$_{1-x}$(BaBi$_2$Ti$_4$O$_{15}$)$_x$, solid solutions ($m = 4$).
indicated that the end member Na$_{0.5}$Bi$_{4.5}$Ti$_4$O$_{15}$ and BaBi$_4$Ti$_4$O$_{15}$ exhibited typical ferroelectric and RFE behaviors, respectively. While, relaxorlike behavior appears for $x/0.6$. But the ferroelectric and piezoelectric properties of Na$_{0.5}$Bi$_{4.5}$Ti$_4$O$_{15}$ x solid solutions were not investigated.

As members of BLSFs with $m=2$, CaBi$_2$Nb$_2$O$_9$ (CBNO) and BaBi$_2$Nb$_2$O$_9$ (BBNO) are particularly interesting because they represent two extreme cases considering the average ionic radii of their A-site cations $r_{Ba}^{2+}=1.61$ Å, $r_{Ca}^{2+}=1.34$ Å, coordination number=12 (Ref. 22) and their dielectric behaviors. CBNO and BBNO exhibit classic ferroelectric$^{23}$ and relaxor$^{24}$ behavior, respectively. Recently, we reported that CBNO has the highest-known Curie point (943 °C) among BLSFs and the highest thermal depoling temperature for polycrystalline ferroelectric ceramics.$^{23,25}$ This means that CBNO is a very promising material for high-temperature (up to 800 °C) piezoelectric applications. How to further improve the piezoelectric activity of CBNO ceramic without dramatically decreasing its high $T_c$ is of great interest for commercial applications. We have also demonstrated that the polarization-electrical field ($P–E$) response of BBNO ceramics is dominated by lossy capacitor behavior and the piezoelectric constant $d_{33}$ is zero after poling at room temperature.$^{24}$ The objective of the current research was to investigate for CBNO–BBNO solid solution if: (1) piezoelectric properties of CBNO can be improved while maintaining their high $T_c$ and (2) RFE behavior occurs with $T_m>300$ °C. In the present study, we present a systematic study of the microstructure and electrical properties of solid solutions in the CBNO–BBNO series with above two aims.

### II. EXPERIMENTAL

Polycrystalline ceramic samples with the formula of (CBNO)$_{1-x}$(BBNO)$_x$, where $x=0$, 0.1, 0.3, 0.5, 0.8, 0.9, and 1.0 were prepared by conventional solid-state reaction sintering. The starting materials were CaCO$_3$ of 99.0% purity, BaCO$_3$ of 99.0% purity, Bi$_2$O$_3$ of 99.975% purity, and Nb$_2$O$_5$ of 99.9% purity. The powders were ball milled for 12

<table>
<thead>
<tr>
<th>$x$</th>
<th>Sintering temperature (°C)</th>
<th>Density (g/cm$^3$)</th>
</tr>
</thead>
<tbody>
<tr>
<td>0</td>
<td>1150</td>
<td>6.82 (96.7%)</td>
</tr>
<tr>
<td>0.1</td>
<td>1150</td>
<td>6.87</td>
</tr>
<tr>
<td>0.3</td>
<td>1150</td>
<td>7.16</td>
</tr>
<tr>
<td>0.5</td>
<td>1100</td>
<td>6.90</td>
</tr>
<tr>
<td>0.8</td>
<td>1100</td>
<td>7.14</td>
</tr>
<tr>
<td>0.9</td>
<td>1100</td>
<td>7.27</td>
</tr>
<tr>
<td>1.0</td>
<td>1100</td>
<td>7.17 (96.9%)</td>
</tr>
</tbody>
</table>

**FIG. 1.** XRD patterns of (CBNO)$_{1-x}$(BBNO)$_x$ ceramics as a function of $x$: (a) $2\theta=10^\circ–80^\circ$ and (b) $2\theta=27^\circ–34^\circ$.

**FIG. 2.** Secondary electron SEM micrographs showing typical microstructures of (CBNO)$_{1-x}$(BBNO)$_x$ ceramics: (a) $x=0$; (b) $x=0.5$; and (c) $x=1.0$.
h in a nylon pot and then calcined in air at 950 °C for 4 h. The calcined powders were remilled for 24 h and admixed with about 3.0 wt % polyethylene glycol as a binder. Subsequently, the powder were dried, sieved to under 500 μm, and uniaxially pressed into disks with 10 mm diameter in a steel die at a pressure of about 200 MPa. The pressed samples were sintered at 1100–1150 °C for 1 h in air. Their density was measured by the Archimedes immersion method.

TABLE II. The fitting results of relaxor compositions for V–F law [Eq. (1)] for the (CBNO)_{1−x}(BBNO)_{x} ceramics.

<table>
<thead>
<tr>
<th>x</th>
<th>0.8</th>
<th>0.9</th>
<th>1.0</th>
</tr>
</thead>
<tbody>
<tr>
<td>$E_a$ (eV)</td>
<td>0.11</td>
<td>0.11</td>
<td>0.56</td>
</tr>
<tr>
<td>$\omega_r$ (Hz)</td>
<td>$1.54 \times 10^{11}$</td>
<td>$5.78 \times 10^{11}$</td>
<td>$4.71 \times 10^{13}$</td>
</tr>
<tr>
<td>$T_f$ (°C)</td>
<td>252</td>
<td>145</td>
<td>−173</td>
</tr>
<tr>
<td>$R^2$</td>
<td>0.9836</td>
<td>0.9862</td>
<td>0.9968</td>
</tr>
</tbody>
</table>

FIG. 3. (Color online) Temperature dependence of dielectric constant $\varepsilon_r$ and loss factor $D$ at various x values in (CBNO)_{1−x}(BBNO)_{x} with different frequencies: (a) x=0; (b) x=0.1; (c) x=0.3; (d) x=0.5; (e) x=0.8; (f) x =0.9; and (g) x=1.0.

FIG. 4. (Color online) Temperature dependence of dielectric constant of (CBNO)_{1−x}(BBNO)_{x} ceramics at the frequency of 1 MHz.
For each composition, there existed a temperature at which the densest ceramics were obtained. Ceramics sintered at temperatures were chosen for final microstructure and property characterizations. Electrodes for electrical property measurements were fabricated with platinum paste (Gwent Electronic Materials Ltd., C2011004D5). The Pt electrodes were fired at 900 °C for 30 min.

X-ray diffraction (XRD) using Cu $K_{\alpha}$ radiation was performed at room temperature for crushed ceramics with a Siemens D5000 to determine the formation of the desired Aurivillius phase. The microstructures of the ceramic specimens were analyzed with a scanning electron microscope (SEM; JEOL JSM 6300). The samples for the SEM study were polished and then thermally etched at about 70 °C below their sintering temperatures for 20 min.

The temperature dependence of the dielectric constants and losses were measured at different frequencies using an LCR meter connected with high-temperature furnace and losses were measured at different frequencies using an LCR meter connected with high-temperature furnace and losses were measured at different frequencies using an LCR meter connected with high-temperature furnace. Since the $T_c$ of CBNO ceramics that consisted of a single Aurivillius phase over the entire composition range without detectable secondary phases within the sensitivity of the XRD equipment. The (020) and (200) peaks of the orthorhombic structure gradually became a single peak with increasing BBNO content $x$, which is expected when Ca$^{2+}$ ions are substituted by larger Ba$^{2+}$ ions on the A-sites. The similar evolution of crystal structure has been reported in solid solution of Bi$_3$TiNbO$_9$–BaBi$_2$Nb$_2$O$_9$.

Figure 2 illustrates typical SEM images of (CBNO)$_{1-x}$(BBNO)$_x$ ceramics. In all cases a randomly oriented platelike morphology developed. This platelike morphology is typical of Aurivillius ceramics and is due to the anisotropic nature of their crystal structure.

### III. RESULTS AND DISCUSSION

Table I summarizes the optimal sintering temperatures and the densities of the (CBNO)$_{1-x}$(BBNO)$_x$ ceramics that were used for final microstructure and property characterizations.

#### A. Crystal structure and morphology

XRD patterns of crushed sintered samples of (CBNO)$_{1-x}$(BBNO)$_x$ are presented in Figs. 1(a) and 1(b). The diffraction patterns were indexed according to the diffraction data of pure CBNO (Joint Committee on Powder Diffraction Standards No. 49-608). It can be seen from Fig. 1(a) that all the patterns are similar and the peaks at (115), (020), and (200) are shifted toward lower 2θ values with increasing BBNO content $x$ [Fig. 1(b)]. The ceramics consisted of a single Aurivillius phase over the entire composition range without detectable secondary phases within the sensitivity of the XRD equipment. The (020) and (200) peaks of the orthorhombic structure gradually became a single peak with increasing BBNO content $x$, which is expected when Ca$^{2+}$ ions are substituted by larger Ba$^{2+}$ ions on the A-sites. The similar evolution of crystal structure has been reported in solid solution of Bi$_3$TiNbO$_9$–BaBi$_2$Nb$_2$O$_9$.

Figure 2 illustrates typical SEM images of (CBNO)$_{1-x}$(BBNO)$_x$ ceramics. In all cases a randomly oriented platelike morphology developed. This platelike morphology is typical of Aurivillius ceramics and is due to the anisotropic nature of their crystal structure.

#### B. Dielectric properties

The dielectric constant $\varepsilon_r$ and loss factor $D$ were measured as a function of temperature at different frequencies ranging from 1 kHz to 1 MHz (Fig. 3). For CBNO ($x=0$) at high frequencies (100 kHz and 1 MHz), a ferroelectric to paraelectric phase transition was apparent at 940 °C, as shown by the sharp frequency independent maximum of $\varepsilon_r$ and the associated minimum in loss [Fig. 3(a)]. Loss peaks observed in CBNO ($x=0$) at a few degrees below the Curie point $T_c$ [Fig. 3(a)], are ascribed to the motion of domain walls. The losses then decreased as a result of the disappearance of the ferroelectricity. Above $T_c$ the losses increased dramatically because of increasing conductivity.

The dielectric behavior for composition with $x=0.1$ [Fig. 3(b)] are similar to that of CBNO ($x=0$). There was a sharp dielectric constant peak at high frequencies ($f \approx 100$ kHz) at $T_c$.

The samples with $x=0.3$ and 0.5 exhibited broad fre-

<table>
<thead>
<tr>
<th>$x$</th>
<th>$0$</th>
<th>$0.1$</th>
<th>$0.3$</th>
<th>$0.5$</th>
<th>$0.8$</th>
<th>$0.9$</th>
<th>$1.0$</th>
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<tbody>
<tr>
<td>$C_1$ (K)</td>
<td>$3.47 \times 10^4$</td>
<td>$4.84 \times 10^4$</td>
<td>$1.44 \times 10^5$</td>
<td>$1.34 \times 10^5$</td>
<td>$\ldots$</td>
<td>$\ldots$</td>
<td>$\ldots$</td>
</tr>
<tr>
<td>$C_2$ (K)</td>
<td>$\ldots$</td>
<td>$\ldots$</td>
<td>$\ldots$</td>
<td>$9.33 \times 10^6$</td>
<td>$1.54 \times 10^7$</td>
<td>$1.21 \times 10^7$</td>
<td></td>
</tr>
<tr>
<td>$T_c$ (°C)</td>
<td>$940$</td>
<td>$918$</td>
<td>$730$</td>
<td>$591$</td>
<td>$\ldots$</td>
<td>$\ldots$</td>
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<tr>
<td>$T_m$ (°C)</td>
<td>$\ldots$</td>
<td>$\ldots$</td>
<td>$\ldots$</td>
<td>$322$</td>
<td>$260$</td>
<td>$241$</td>
<td></td>
</tr>
<tr>
<td>$\theta$ (°C)</td>
<td>$845$</td>
<td>$824$</td>
<td>$284$</td>
<td>$149$</td>
<td>$\ldots$</td>
<td>$\ldots$</td>
<td></td>
</tr>
<tr>
<td>$\gamma$</td>
<td>$\ldots$</td>
<td>$\ldots$</td>
<td>$\ldots$</td>
<td>$1.73$</td>
<td>$1.82$</td>
<td>$1.94$</td>
<td></td>
</tr>
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</table>
quency independent dielectric constant peaks, as shown in Figs. 3(c) and 3(d). A similar dielectric behavior was also found for SrBi$_4$Ti$_4$O$_{15}$ doped by barium.$^{29}$

The dielectric behavior of the compositions with $x=0.8$, 0.9, and 1.0 [Figs. 3(e)–3(g)] were very different from that of the compositions with smaller $x$ values ($x \leq 0.5$). The following features were observed: (i) the phase transition was no longer sharp but diffuse; (ii) the temperature $T_m$ of the maximum $\varepsilon_r$ increased with increasing frequency and the magnitude of the $\varepsilon_r$ decreased; and (iii) the losses exhibited a maximum at temperatures below $T_m$ and this loss maximum was also shifted toward higher temperatures as the frequency increased. All these features are characteristic of ferroelectric relaxor behavior.$^{30}$ It is interesting to note that $T_m$ of composition with $x=0.8$ shifts from 308 °C at 10 kHz to 322 °C at 1 MHz, which indicates it is a RFE composition with $T_m$ temperatures higher than that of lead-based perovskite materials.

Since the frequency dispersion of relaxors does not follow the Debye law, the Vogel–Fulcher (V–F) law is used to describe the shift in the dielectric peaks with frequency$^{31}$

$$\omega = \omega_0 \exp \left( - \frac{E_a}{k(T - T_f)} \right), \quad (1)$$

where $\omega_0$ is the attempt frequency, which is associated with the cut-off frequency of the distribution of relaxation times. $\omega$ is the measurement frequency. $E_a$ is the energy barrier for dipole switching and $T_f$ is a freezing temperature, at which all relaxation times diverge and the distribution of relaxation times, $\tau$, becomes infinitely long. The values of fitting parameters for relaxor compositions with $x=0.8$, 0.9, and 1.0 are shown in Table II. The fitted $\omega_0$ values are in the frequency range of lattice vibration for ionic solids, which is of the order of $10^{11}–10^{13}$ Hz.$^{32}$ The fitted $T_f$ temperature was 252 °C for $x=0.8$, 145 °C for $x=0.9$, and −173 °C for $x=1.0$, respectively.

Figure 4 shows the dielectric constants of (CBNO)$_{1-x}$ (BBNO)$_x$ ceramics as a function of temperature measured at 1 MHz. It clearly shows that with increasing

FIG. 6. (Color online) $I$–$E$ (blue line) and $P$–$E$ (red line) loops of (CBNO)$_{1-x}$ (BBNO)$_x$ ceramics measured at 200 °C and 10 Hz.
barium content \( x \) on the A-sites the temperature of maximum \( \varepsilon_r \) decreased and dielectric peak became broader, which is consistent with previous studies of the effect of Ba-doping on the dielectric properties of Aurivillius materials.\textsuperscript{21,33–36}

For compositions with \( x=0–0.5 \), the fitting parameters for the Curie–Weiss law \( \varepsilon^{-1}=(T-\theta)/C_1 \) are given in Table III. As expected for BLSFs, the Curie constants \( C_1 \) calculated from the values of the permittivity above \( T_c \) at 1 MHz are of the order of \( 10^4–10^5 \) K,\textsuperscript{37} which suggests that \( (\text{CBNO})_{1-x}(\text{BBNO})_x \) with \( x=0–0.5 \) are typical displacive-type ferroelectrics.\textsuperscript{36} This is consistent with the studies on the origin of ferroelectricity in BLSFs, as mentioned early. The extrapolated Curie–Weiss temperatures \( \theta \) are well below the Curie points, which can be attributed to the high losses near the Curie points. The dielectric behaviors of the compositions with \( x \geq 0.8 \) show a diffuse phase transition and depart from the classic Curie–Weiss law. A modified Curie–Weiss law \( \varepsilon^{-1}-\varepsilon_m=(T-T_m)^\gamma/C_2 \) \( (T_m \) is the temperature of the maximum permittivity \( \varepsilon_m \) at a given frequency, \( \gamma=1 \) for normal ferroelectrics and \( \gamma=2 \) for ideal relaxor) is often used to fit the permittivity data.\textsuperscript{39} The \( \gamma \) values obtained at 1 MHz are close to 2, confirming the relaxor nature of these compositions (Table III). The Curie constants \( C_2 \) for \( x \geq 0.8 \) are one or two orders of magnitude higher than the \( C_1 \) values for the compositions with \( x=0–0.5 \), which is consistent with the previous studies on relaxor behavior in BLSFs compounds.\textsuperscript{21,36}

It is well known that the Curie point is closely related to structural displacements.\textsuperscript{27} The degree of distortion of the perovskite structure can be determined by the tolerance factor \( t, t=(r_A+r_O)/\sqrt{2}(r_B+r_O) \), where \( r_A, r_B, \) and \( r_O \) are the ionic radii of the A and B-site cations, and an oxygen ion, respectively.\textsuperscript{27} \( r_{Ca}^+=1.34 \) Å, \( r_{Ba}^+=1.61 \) Å, \( r_{A}^+=1.40 \) Å, \( r_{B}^+=0.64 \) Å, \( r_{O}^-=1.40 \) Å, \( CN=12; r_{Ba}^+=1.61 \) Å, \( CN=6; r_{O}^-=1.40 \) Å, \( CN=6).\textsuperscript{22} Figure 5 shows the relationship between the tolerance factor \( t \) and \( T_c \) or \( T_m \) of the \( (\text{CBNO})_{1-x}(\text{BBNO})_x \) system at 1 MHz. The tolerance factor \( t \) is found to increase with the increasing barium content \( x \) on the A-sites. This results in a reduction in the tilting of the \( \text{NbO}_6 \) octahedron.\textsuperscript{27} The atomic displacements along the polar axis (a-axis) are smaller so that the energy involved to reach the prototype high-temperature structure is lowered. Accordingly, \( T_c \) or \( T_m \) decreases.

C. Ferroelectric properties

For the \( (\text{CBNO})_{1-x}(\text{BBNO})_x \) ceramics, due to the lower limits of the attainable sample thickness (0.10–0.15 mm) produced by polishing damage, maximum current limit of the high voltage amplifier available (2.5 mA), high conductivities at high-temperature and high coercive fields (\( E_c \)), it was usually not possible to obtain saturated ferroelectric hysteresis loops. Figure 6 shows the \( I-E \) and \( P-E \) loops measured at 10 Hz and 200 °C for all of the compositions studied. The onset of ferroelectric domain switching, as indicated by current peaks in the \( I-E \) curves, was observed for composition with \( x=0.5 \). For compositions with \( x \geq 0.8 \), the currents consisted of the dielectric displacement current and conduction current with no apparent ferroelectric switching peak.\textsuperscript{30} According to the dipolar glass model developed by Viehland et al.\textsuperscript{31} the relaxor behaviors are related to nanodomains. The remanent polarization drops quickly above \( T_f \) because thermal fluctuations destroy the correlation of the nanodomains. The fitting values of \( T_f \) for compositions with \( x=0.8, 0.9, \) and 1.0 are 252, 145, and −173 °C (Table III), respectively. The absence of domain switching in compositions with \( x \geq 0.8 \) is due to the fact that the test temperature (200 °C) is very close to or higher than \( T_f \). The \( I-E \) and \( P-E \) loops for compositions with \( x=0.8 \) were also measured.
at 10 Hz and 25 °C, as shown in Fig. 7. However, no ferroelectric peaks were observed for any of these compositions.

**D. Piezoelectric properties**

Figure 8 shows the dependence of piezoelectric coefficient $d_{33}$ for samples poled at 100 and 200 °C. When the samples were poled at 200 °C, the $d_{33}$ value reached a maximum (10.8 pC/N) for $x=0.1$. The $d_{33}$ value for $x=0.3$ or 0.5 was also higher than that of pure CBNO ($x=0$). The combination of high $d_{33}$ values and high $T_{c}$ points (>700 °C) suggests that the composition with $x=0.3$ could be good candidates for high-temperature piezoelectric applications. Thecomposition with $x=0.8$ showed ferroelectricity, as evidenced by a nonzero $d_{33}$ value of 1.5 pC/N after poling at 100 °C. From the diffuse dielectric behavior [Fig. 3(e)] and the nonzero $d_{33}$, it has been demonstrated that the composition with $x=0.8$ is a RFE with $T_{m} > 300$ °C. Considering that the error of the piezo-$d_{33}$ meter is ±0.2 pC/N, the low $d_{33}$ value (0.3 pC/N) for the composition with $x=0.9$ suggests very weak ferroelectricity after poling at 100 °C. The existence of nonzero $d_{33}$ for compositions with $x=0.8$ and 0.9 after poling at 100 °C is consistent with the fact that the fitted $T_{c}$ values of these two compositions (see Table II) are higher than 100 °C. The $d_{33}$ value of BBNO ($x=1.0$) was zero when poled at 100 and 200 °C, even though the samples were poled under high field of 150 kV/cm. Ando et al.10 also reported that no piezoelectricity was observed in BBNO.

**IV. CONCLUSION**

(CBNO)$_{1-x}$(BBNO)$_{x}$ solid solution ceramics were fabricated by conventional solid-state reaction. The substitution of Ca$^{2+}$ with the larger Ba$^{2+}$ cation on the A-sites of the perovskite slabs produced a decrease in $T_{c}$ or $T_{m}$ and a broadening of the temperature maximum of the dielectric permittivity. Relaxor behavior occurred at compositions with $x \leq 0.8$. Compositions with $x \leq 0.5$ exhibited ferroelectric behavior, as evidenced by ferroelectric domain switching peaks in $I$-$E$ loops measured at 200 °C. The $d_{33}$ value of CBNO ceramics can be effectively enhanced by Ba$^{2+}$ doping on A-sites ($x \leq 0.3$). The combination of high $d_{33}$ values and high $T_{c}$ points (>700 °C) suggests that the compositions with $x=0.3$ could be good candidates for high-temperature piezoelectric application. The composition with $x=0.8$ was a relaxor ferroelectric with $T_{m}$ around 320 °C at 1 MHz.